

# Formation of interstellar methanol ice prior to the heavy CO freeze-out stage

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## ABSTRACT

**Context.** The formation of methanol (CH<sub>3</sub>OH) on icy grain mantles during the star formation cycle is mainly associated with the CO freeze-out stage. Yet there are reasons to believe that CH<sub>3</sub>OH also can form at an earlier period of interstellar ice evolution in CO-poor and H<sub>2</sub>O-rich ices.

**Aims.** This work focuses on CH<sub>3</sub>OH formation in a H<sub>2</sub>O-rich interstellar ice environment following the OH-mediated H-abstraction in the reaction, CH<sub>4</sub> + OH. Experimental conditions are systematically varied to constrain the CH<sub>3</sub>OH formation yield at astronomically relevant temperatures.

**Methods.** CH<sub>4</sub>, O<sub>2</sub>, and hydrogen atoms are co-deposited in an ultrahigh vacuum chamber at 10–20 K. OH radicals are generated by the H + O<sub>2</sub> surface reaction. Temperature programmed desorption – quadrupole mass spectrometry (TPD–QMS) is used to characterize CH<sub>3</sub>OH formation, and is complemented with reflection absorption infrared spectroscopy (RAIRS) for CH<sub>3</sub>OH characterization and quantitation.

**Results.** CH<sub>3</sub>OH formation is shown to be possible by the sequential surface reaction chain, CH<sub>4</sub> + OH → CH<sub>3</sub> + H<sub>2</sub>O and CH<sub>3</sub> + OH → CH<sub>3</sub>OH at 10–20 K. This reaction is enhanced by tunneling, as noted in a recent theoretical investigation Lamberts et al. (2017, A&A, 599, A132). The CH<sub>3</sub>OH formation yield via the CH<sub>4</sub> + OH route versus the CO + H route is approximately 20 times smaller for the laboratory settings studied. The astronomical relevance of the new formation channel investigated here is discussed.

**Key words.** astrochemistry

## 1. Introduction

Methanol (CH<sub>3</sub>OH) is an important interstellar molecule. CH<sub>3</sub>OH has been observed abundantly in both the gas phase (Friberg et al. 1988; Turner 1998; Parise et al. 2002; Bergman et al. 2011; Wirström et al. 2011; Guzmán et al. 2013; Öberg et al. 2014; Taquet et al. 2015) and the solid state (Grim et al. 1991; Allamandola et al. 1992; Skinner et al. 1992; Chiar et al. 1996; Dartois et al. 1999; Gibb et al. 2000, 2004; Pontoppidan et al. 2003; Taban et al. 2003; Boogert et al. 2008, 2011; Bottinelli et al. 2010). It is generally accepted that CH<sub>3</sub>OH formation is most efficient by solid state interactions on icy grain mantles. Models, supported by experiments, of gas phase synthesis of CH<sub>3</sub>OH provide abundances orders of magnitude below the observed fractional abundance of CH<sub>3</sub>OH (Garrod et al. 2006; Geppert et al. 2006), while solid state laboratory studies show that CH<sub>3</sub>OH is efficiently formed in CO-rich ices through sequential hydrogenation of CO (Hiraoka et al. 1994; Watanabe & Kouchi 2002; Fuchs et al. 2009). This is further supported by computational models (Cuppen et al. 2009; Fuchs et al. 2009; Garrod et al. 2006) that show that hydrogenation of CO ice leads to the production of CH<sub>3</sub>OH. Indeed, these findings are in line with the spectroscopic interpretation of observational data; CO and CH<sub>3</sub>OH are found to coexist in CO-rich and

H<sub>2</sub>O-poor interstellar ices (Cuppen et al. 2011; Boogert et al. 2015; Penteado et al. 2015). Not only is it a prevalent interstellar molecule, but CH<sub>3</sub>OH is also an important precursor in the formation of larger species. Öberg et al. (2009) illustrated that upon vacuum UV irradiation, CH<sub>3</sub>OH can break apart into fragments that can recombine to form complex organic molecules (COMs). In UV-rich environments in the interstellar medium, the formation of CH<sub>3</sub>OH may thus be crucial to the formation of larger COMs. The studies by Chuang et al. (2015, 2017) and Fedoseev et al. (2017) show that CH<sub>3</sub>OH may be a key player in the formation of larger COMs, also for cold dense prestellar core conditions. The authors show that radicals derived from CH<sub>3</sub>OH via an abstraction process can recombine or combine with other radicals, resulting in the formation of COMs as large as glycerol, even at temperatures below 20 K and without the need for UV-induced radiation.

Recent theoretical and experimental efforts (Lamberts et al. 2017) have provided results that have sparked the idea of H<sub>2</sub>O and CH<sub>3</sub>OH coexisting in a CO-poor and H<sub>2</sub>O-rich interstellar ice. Such ices are thought to form before the “heavy” CO freeze out stage, which starts to occur at a cloud extinction ( $A_V$ ) of  $A_V > 3$  and dust temperatures  $< 20$  K, and that is followed by the “catastrophic” CO freeze-out stage at  $A_V > 9$  and dust temperatures also  $< 20$  K. Prior to the heavy CO freeze-out stage,

a CO:H<sub>2</sub>O ice ratio of <5% is expected due to some CO freezing out as well as atom-addition reactions (C, O, H, etc.). Due to the relatively low gas densities at this stage, only a relatively low amount of CO is able to accrete onto dust grains, as examined in Pontoppidan (2006) and discussed in Öberg et al. (2011), Boogert et al. (2015), and van Dishoeck (2017). This time period is known as the “H<sub>2</sub>O-rich ice phase” or “polar ice phase” (i.e., the phase that lacks CO ice). CH<sub>3</sub>OH formation through this reaction, CH<sub>4</sub> + OH, can take place in this phase. However, this does not exclude this reaction to also take place during the CO freeze-out stage, where CH<sub>3</sub>OH formation is normally dominated by CO hydrogenation. During this time, O<sub>2</sub> might be intimately mixed with CO (Vandenbussche et al. 1999) and can be hydrogenated to form OH radicals. Gas phase CH<sub>4</sub> becomes more dominant relative to atomic carbon (van Dishoeck 1998) and could freeze out with CO. The importance of the present study is that it focuses on a phase in which CH<sub>3</sub>OH is only formed through CH<sub>4</sub> + OH, as CO is not frozen out yet, and an identification of CH<sub>3</sub>OH in corresponding environments, therefore, directly relates to this alternative reaction pathway.

The study by Lamberts et al. (2017) was done to provide reaction rates at cryogenic temperatures of the reaction between methane (CH<sub>4</sub>) and OH as CH<sub>4</sub> has been observed in the H<sub>2</sub>O-rich ice phase (Boogert et al. 1997; Öberg et al. 2008), and OH radicals are expected to be abundant in that phase (Cuppen & Herbst 2007; Ioppolo et al. 2008; Cuppen et al. 2010; Oba et al. 2012; Lamberts et al. 2013). These authors, as well as others (Wada et al. 2006; Hodyss et al. 2009; Weber et al. 2009; Zins et al. 2012; Bossa et al. 2015), showed that CH<sub>4</sub> hydrogen abstraction by an OH radical results in the efficient formation of CH<sub>3</sub> radicals; a process that can be induced by tunneling (Lamberts et al. 2017). Thus, these findings can be taken one step further by postulation that the OH-induced abstraction reaction of CH<sub>4</sub> may result in the formation of CH<sub>3</sub>OH by the sequential reaction chains: CH<sub>4</sub> + OH → CH<sub>3</sub> + H<sub>2</sub>O and CH<sub>3</sub> + OH → CH<sub>3</sub>OH. There are already a number of laboratory experiments that exhibit CH<sub>3</sub>OH formation via energetic processing of ice mixtures containing CH<sub>4</sub> and H<sub>2</sub>O (Schutte 1988; d’Hendecourt et al. 1996; Moore & Hudson 1998; Wada et al. 2006; Hodyss et al. 2009; Martín-Doménech et al. 2016). In addition, Bergner et al. (2017) have explored the formation of CH<sub>3</sub>OH via O-atom insertion in CH<sub>4</sub> molecules upon photo dissociation of O<sub>2</sub> in binary mixtures. In our paper, the reaction CH<sub>4</sub> + OH leading to the formation of CH<sub>3</sub>OH under relevant cold dense core conditions is discussed. The OH radicals are already present at the earliest stages of ice formation via photo dissociation of H<sub>2</sub>O molecules and by atomic H- and O-rich accreting gas (see, e.g., Boogert et al. 2015). Therefore, this formation mechanism could work even under non-energetic conditions, and can be extended to the formation of COMs that contain the hydroxyl functional group (–OH). Furthermore, CH<sub>3</sub> radical formation is constrained to abstraction reactions only, whereas in the energetic processing studies, it is unclear whether CH<sub>3</sub> radicals are formed by energetic and/or non-energetic pathways.

Data from astronomical observations also provide an incentive to explore CH<sub>3</sub>OH ice formation before the heavy CO freeze-out stage (i.e., in the H<sub>2</sub>O-rich ice phase). In the review published by Boogert et al. (2015), Fig. 7 displays the observed ice column densities of CH<sub>3</sub>OH and H<sub>2</sub>O, as well as other species, as a function of ( $A_V$ ). Consistent with the current model of interstellar ice evolution, H<sub>2</sub>O is concentrated at a lower  $A_V$  with respect to CH<sub>3</sub>OH. However, taking a closer look at the figure reveals that CH<sub>3</sub>OH upper limit values can be found at a lower  $A_V$  and at lower column densities than CH<sub>3</sub>OH detections.

Such upper limits do not prove that CH<sub>3</sub>OH can be formed at lower  $A_V$ , but they also do not exclude the idea. A possible explanation for the relatively low  $A_V$  and column densities for the CH<sub>3</sub>OH upper limits may be due to other reaction pathways expected to take place before the CO freeze-out phase. It is also possible that at such low  $A_V$ , the sensitivity of the telescopes used may not have been high enough. In the dawn of the *James Webb Space Telescope* (JWST) era, it is expected that soon it will be possible to extend on the formation schemes in which solid CH<sub>3</sub>OH is formed or consumed.

In this paper, the formation of CH<sub>3</sub>OH ice by CH<sub>4</sub> + OH under relevant cold dense core conditions is discussed. Section 2 provides an overview of the experimental conditions and methods. Section 3 details how CH<sub>3</sub>OH was identified, formed in the ice, and quantified. Section 4 provides insights into how the laboratory results presented here can be used to constrain the ice chemistry in the H<sub>2</sub>O-rich ice phase of cold prestellar core environments, as well as a discussion of Fig. 7 from Boogert et al. (2015). Section 5 summarizes the findings from this paper.

## 2. Experimental setup and methods

The experiments described here are performed with SURFRE-SIDE<sup>2</sup>. This is a double atom beam line, ultrahigh vacuum (UHV) setup with a base pressure of  $\sim 10^{-10}$  mbar in the main chamber. An in-depth description of the setup and experimental procedures can be found in Ioppolo et al. (2013) and Linnartz et al. (2015). Ices are grown on a gold-plated copper substrate ( $2.5 \times 2.5$  cm<sup>2</sup>) that is attached to the cold-finger of a closed-cycle helium cooled cryostat (ColdEdge CH-204 SF). Temperatures as low as 7 K can be realized. The temperature settings are controlled by a LakeShore 340 temperature controller. The substrate temperature is measured by a silicon diode sensor (DT-670) with an absolute accuracy of 1 K. The substrate temperature is changed by resistive heating of a tungsten filament. Incorporation of a sapphire rod within the cryocooler allows substrate temperatures as high as 450 K.

The processes taking place in the ice can be studied through infrared (IR) spectroscopy and mass spectrometry. Species formed in the ice are probed by their IR signatures with a Fourier transform infrared spectrometer (FTIR; Agilent Cary 640/660) applying the reflection absorption infrared spectroscopy (RAIRS) technique. The FTIR permits a coverage of 6000–700 cm<sup>-1</sup> with a resolution of 1 cm<sup>-1</sup>. To further constrain species present in the ice, a quadrupole mass spectrometer (QMS; Spectra Microvision Plus LM76) is used to measure the desorption temperature of ice species during a temperature programmed desorption (TPD) run, as well as the mass spectrum of the desorbing species upon electron impact ionization. A commonly used electron impact ionization energy of 70 eV is chosen. All TPD experiments involve a linear ramp rate of 2 K min<sup>-1</sup>.

In order to simulate cold dense core conditions, all ices are grown on a substrate surface with a temperature between 10–20 K. CH<sub>4</sub>, H, and O<sub>2</sub> are admitted into the main vacuum chamber following the co deposition technique (i.e., different molecular species deposited simultaneously), which reproduces interstellar conditions better than with previously used pre-deposition techniques (i.e., different molecular species deposited sequentially; Linnartz et al. 2015). Additionally, this technique allows all deposited species to react with one another, regardless of the ice thickness, which is advantageous when trying to probe trace species. The OH radicals are formed by H-atom addition to O<sub>2</sub> (Cuppen et al. 2010) and the CH<sub>3</sub> radicals are formed by

Table 1. Experiments performed in this study.

No.	Experiments	Ratio CH <sub>4</sub> :H:O <sub>2</sub>	T <sub>sample</sub> <sup>a</sup> (K)	Flux <sub>CH<sub>4</sub></sub> cm <sup>-2</sup> s <sup>-1</sup>	Flux <sub>H</sub> cm <sup>-2</sup> s <sup>-1</sup>	Flux <sub>O<sub>2</sub></sub> cm <sup>-2</sup> s <sup>-1</sup>	Flux <sub>CH<sub>3</sub>OH</sub> cm <sup>-2</sup> s <sup>-1</sup>	Flux <sub>H<sub>2</sub>O</sub> cm <sup>-2</sup> s <sup>-1</sup>	Flux <sub>CO</sub> cm <sup>-2</sup> s <sup>-1</sup>	Time (s)
1.0	CH <sub>4</sub> + H + O <sub>2</sub>	1:2:1	10	3E12	6E12	4E12	–	–	–	43 200
2.0	CH <sub>4</sub> + H + O <sub>2</sub>	1:2:1	10	3E12	6E12	4E12	–	–	–	21 600
2.1	<sup>13</sup> CH <sub>4</sub> + H + O <sub>2</sub>	1:2:1	10	3E12	6E12	4E12	–	–	–	21 600
2.2	CH <sub>4</sub> + H + <sup>18</sup> O <sub>2</sub>	1:2:1	10	3E12	6E12	4E12	–	–	–	21 600
2.3	<sup>13</sup> CH <sub>4</sub> + H + <sup>18</sup> O <sub>2</sub>	1:2:1	10	3E12	6E12	4E12	–	–	–	21 600
3.0	CH <sub>3</sub> OH	–	10	–	–	–	4E13	–	–	100
4.0	CH <sub>3</sub> OH	–	10	–	–	–	1E13	–	–	1200
4.1	CH <sub>3</sub> OH + O <sub>2</sub> (MWAS) <sup>b</sup>	–	10	–	–	4E13	–	–	–	9000
4.2	CH <sub>3</sub> OH + CH <sub>4</sub>	–	10	6E13	–	–	–	–	–	9000
4.3	CH <sub>3</sub> OH + H <sub>2</sub> O	–	10	–	–	–	–	6E13	–	9000
4.4	CH <sub>3</sub> OH + H <sub>2</sub> O (MWAS) <sup>b</sup>	–	10	–	–	–	–	6E13	–	9000
4.5	CH <sub>3</sub> OH + H <sub>2</sub> O (MWAS) <sup>b</sup> + CH <sub>4</sub>	–	10	–	–	–	–	6E13	–	9000
4.6	CH <sub>3</sub> OH + CO	–	15	6E13	–	–	–	6E13	–	9000
5.0	<sup>13</sup> CH <sub>4</sub> + H + <sup>18</sup> O <sub>2</sub>	1:2:1 <sup>c</sup>	20	3E12 <sup>c</sup>	–	–	2E10	–	7E10	3600
6.0	CH <sub>4</sub> + H + O <sub>2</sub>	1:2:1	10	3E12	6E12	4E12	–	–	–	21 600
					6E12	1E13	–	–	–	21 600

Notes. All fluxes, except the H flux, are derived from the Hertz–Knudsen equation. <sup>(a)</sup> Temperature of the sample at which the ices are grown. <sup>(b)</sup> The CH<sub>4</sub> flux is lower than the listed value in order to have the same CH<sub>4</sub>:H<sub>2</sub>O ice ratio as found in exp. 2.0. Thus, a constant CH<sub>4</sub> flux is not carried out in this particular experiment. <sup>(c)</sup> “(MWAS)” denotes species that were placed in the MWAS chamber. See also Sect. 3.2.

OH-mediated H-abstraction of CH<sub>4</sub> (i.e., radicals are formed in situ). A series of control experiments performed by Lamberts et al. (2017) showed that the produced OH radicals are solely responsible for the H-abstraction of CH<sub>4</sub> (i.e., formation of CH<sub>3</sub> radicals directly through reaction with H atoms is not efficient, as proved in Lamberts et al. (2017)). We note that the formation of OH radicals from O<sub>2</sub> is not necessarily representative of the formation of OH radicals in low A<sub>V</sub> (low density) environments, where the O + H route dominates, as well as photodissociation of H<sub>2</sub>O. Since the OH and CH<sub>3</sub> radicals are formed in the ice, they are expected to be thermalized before further reactions occur.

To perform H-atom addition reactions, H-atoms are produced by a hydrogen atom beam source (HABS; Tschersich & Von Bonin 1998; Tschersich 2000; Tschersich et al. 2008), and the H-atom beam line has an angle of 45° to surface normal of the gold-plated sample. Hydrogen molecules (H<sub>2</sub>; Praxair 99.8%) flow into the HABS chamber via a leak valve and are thermally cracked by a tungsten filament. The gas deposition lines of O<sub>2</sub> (Linde Gas 99.999%) and CH<sub>4</sub> (Linde Gas 99.995%) are angled at 135° and 68°, respectively, to the plane of the sample’s surface. Gas isotopologues, <sup>18</sup>O<sub>2</sub> (Campro Scientific 97%) and <sup>13</sup>CH<sub>4</sub> (Sigma-Aldrich 99%), are used as controls to aid in identification of the formed ice products.

Fluxes and column densities are characterized as follows. An absolute D-atom flux was measured by Ioppolo et al. (2013), and the O<sub>2</sub> and CH<sub>4</sub> fluxes are calculated using the Hertz–Knudsen equation (Kolasinski 2012). Column densities are determined by the relation between absorbance and column density, as described in Hudgins et al. (1993). As discussed previously by Ioppolo et al. (2013), Teolis et al. (2007), and Loeffler et al. (2006), such measurements must be done with caution. Therefore in this work, only the relative column densities of CH<sub>4</sub>, H<sub>2</sub>O, and CH<sub>3</sub>OH are given (i.e., absolute column densities are not listed). Care is taken to use integrated absorbances that do not deviate from the linear trend of the column density over time. The band strength used to determine the CH<sub>3</sub>OH column density is  $7.1 \times 10^{-17}$  cm molecule<sup>-1</sup> (1030 cm<sup>-1</sup>) and is obtained by performance of a He–Ne laser interference experiment in SURFRESIDE<sup>2</sup> (Chuang et al. in prep.). The underlying experimental procedure is described in detail in Paardekooper et al. (2016). CH<sub>4</sub> (1302 cm<sup>-1</sup>) and H<sub>2</sub>O (1659 cm<sup>-1</sup>) band strengths of  $8.0 \times 10^{-18}$  cm molecule<sup>-1</sup> and  $9.0 \times 10^{-18}$  cm molecule<sup>-1</sup>, respectively, are extracted from Bouilloud et al. (2015). Since the band strengths from Bouilloud et al. (2015) are obtained from transmission IR experiments, a proportionality factor between the CH<sub>3</sub>OH band strengths from Bouilloud et al. (2015) and our laser interference experiment is used as a correction factor. It should be noted that literature transmission band strength values cannot be used with optical depth values obtained from a RAIRS experiment partly due to surface-enhanced dipole coupling that occurs in RAIRS. Setup specific values have to be determined for specified conditions. All experiments discussed in this paper are listed in Table 1.

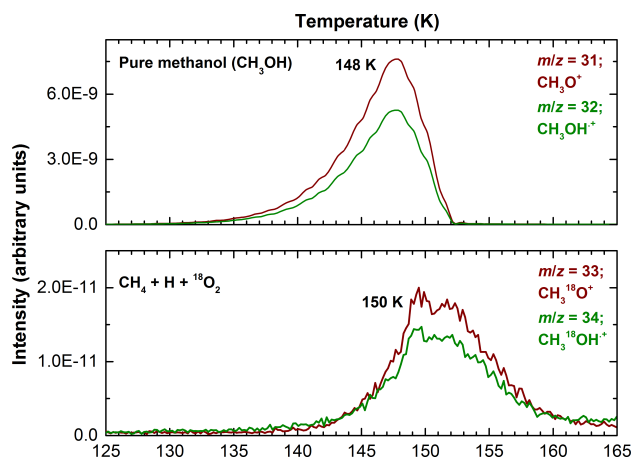
### 3. Results and discussion

#### 3.1. Identification and analysis of CH<sub>3</sub>OH formation

In the next sections, three different ways to confirm the formation of CH<sub>3</sub>OH ice by the reaction, CH<sub>4</sub> + OH, are presented.

##### 3.1.1. TPD

Figure 1 (top) shows the TPD spectra of pure CH<sub>3</sub>OH compared to the TPD spectra of the CH<sub>4</sub> + H + <sup>18</sup>O<sub>2</sub> reaction (bottom)

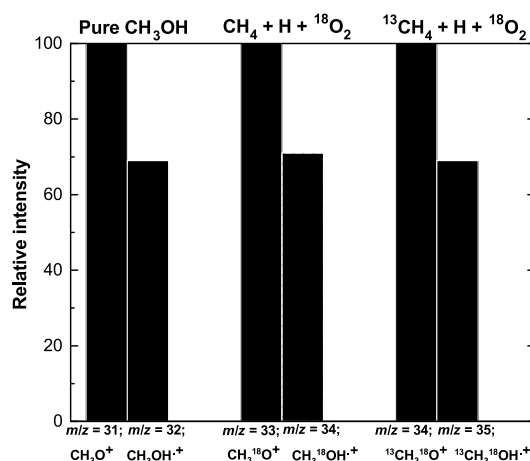


**Fig. 1.** TPD of pure CH<sub>3</sub>OH (*top*; exp. 3.0) and TPD of the CH<sub>4</sub> + H + <sup>18</sup>O<sub>2</sub> reaction (*bottom*; exp. 2.2).

for  $m/z = 31$ ,  $32$  and  $33$ ,  $34$ , respectively. We note that  $m/z = 32$  also represents the O<sub>2</sub><sup>+</sup> ion in the CH<sub>4</sub> + H + O<sub>2</sub> reaction, and therefore TPD data from the <sup>18</sup>O<sub>2</sub> isotope experiment are used instead. Because  $m/z = 31$  (CH<sub>3</sub>O<sup>+</sup>) and  $32$  (CH<sub>3</sub>OH<sup>+</sup>) represent the most intense ions for the CH<sub>3</sub>OH cracking pattern found in our experimental setup and because their isotopically-induced shifted  $m/z$  values in the isotope-enriched experiments can be tracked for most experiments (i.e., the  $m/z$  values remain characteristic of CH<sub>3</sub>OH and not of other species), we use the  $m/z$  values representing CH<sub>3</sub>O<sup>+</sup> and CH<sub>3</sub>OH<sup>+</sup> to confirm the formation of CH<sub>3</sub>OH. As seen in Fig. 1 (*top*), the signals for  $m/z = 31$  and  $32$  peak at 148 K for pure CH<sub>3</sub>OH. The same  $m/z$  signals peak at 150 K in the CH<sub>4</sub> + H + <sup>18</sup>O<sub>2</sub> experiment, as seen in Fig. 1 (*bottom*). This slight increase in the CH<sub>3</sub>OH desorption temperature is expected since the binding energy of CH<sub>3</sub>OH will be influenced by the presence of H<sub>2</sub>O. It should be noted that H<sub>2</sub>O, along with H<sub>2</sub>O<sub>2</sub>, are formed during the experiments and are not initial reactants of the ice mixture. The similar peak positions and profiles for the CH<sub>3</sub>O<sup>+</sup> (CH<sub>3</sub><sup>18</sup>O<sup>+</sup>) and CH<sub>3</sub>OH<sup>+</sup> (CH<sub>3</sub><sup>18</sup>OH<sup>+</sup>) ion signals between the two experiments allows to assign the observed desorption as originating from CH<sub>3</sub>OH.

This assignment is further constrained by quantitatively comparing the CH<sub>3</sub>OH fragmentation pattern upon 70 eV electron impact ionization with values available from the NIST database<sup>1</sup>. Figure 2 presents a column chart of the relative integrated intensities of the CH<sub>3</sub>O<sup>+</sup> and CH<sub>3</sub>OH<sup>+</sup> ion signals observed in a pure CH<sub>3</sub>OH experiment and in isotopically-enriched <sup>13</sup>CH<sub>4</sub> + H + <sup>18</sup>O<sub>2</sub> experiments. As seen in Fig. 2, the relative integrated intensities of pure CH<sub>3</sub>OH for  $m/z = 31$  (CH<sub>3</sub>O<sup>+</sup>) and  $32$  (CH<sub>3</sub>OH<sup>+</sup>) are 100 and 69, respectively. In the CH<sub>4</sub> + H + <sup>18</sup>O<sub>2</sub> experiment, the intensities for  $m/z = 33$  (CH<sub>3</sub><sup>18</sup>O<sup>+</sup>) and  $34$  (CH<sub>3</sub><sup>18</sup>OH<sup>+</sup>) are 100 and 71, respectively. In the <sup>13</sup>CH<sub>4</sub> + H + <sup>18</sup>O<sub>2</sub> experiment, the values for  $m/z = 34$  (<sup>13</sup>CH<sub>3</sub><sup>18</sup>O<sup>+</sup>) and  $35$  (<sup>13</sup>CH<sub>3</sub><sup>18</sup>OH<sup>+</sup>) are 100 and 69, respectively. In the NIST database, the relative integrated intensities for  $m/z = 31$  (CH<sub>3</sub>O<sup>+</sup>) and  $32$  (CH<sub>3</sub>OH<sup>+</sup>) are 100 and 74, respectively. The relative integrated intensities in the three experiments are nearly identical and are close to the values from NIST, fully in line with CH<sub>3</sub>OH formation in the ice.

<sup>1</sup> NIST Mass Spec Data Center, S. E. Stein, director, “Mass Spectra” in NIST Chemistry WebBook, NIST Standard Reference Database Number 69, Eds. P. J. Linstrom and W. G. Mallard, National Institute of Standards and Technology, Gaithersburg MD, 20899, doi: [10.18434/T4D303](https://doi.org/10.18434/T4D303), (retrieved November 7, 2017).



**Fig. 2.** Integrated QMS signals normalized to the integrated QMS signals of the CH<sub>3</sub>O<sup>+</sup>, CH<sub>3</sub><sup>18</sup>O<sup>+</sup>, and <sup>13</sup>CH<sub>3</sub><sup>18</sup>O<sup>+</sup> ions from the CH<sub>3</sub>OH (exp. 3.0), CH<sub>4</sub> + H + <sup>18</sup>O<sub>2</sub> (exp. 2.2), and <sup>13</sup>CH<sub>4</sub> + H + <sup>18</sup>O<sub>2</sub> (exp. 2.3) experiments, respectively. Ratios of the integrated intensities between the isotopes are in agreement with each other and with NIST values representing CH<sub>3</sub>OH.

### 3.1.2. RAIRS

Figure 3 shows the resulting 4000–700 cm<sup>-1</sup> (2.5–14.3 μm) spectrum for CH<sub>4</sub> + H + O<sub>2</sub> interacting on a 10 K substrate. Table 2 lists the frequencies measured, as well as the identification of species observed and the corresponding vibrational modes. The assignment of the peaks are constrained by considering literature values, isotope experiments, ice desorption temperatures, and varying flux ratios of the reactants. The inset in Fig. 3 shows a band at 1001 cm<sup>-1</sup> that is assigned to the C–O stretching mode of CH<sub>3</sub>OH. This assignment is based on experiments performed in Sect. 3.2. Literature values for this band (Wada et al. 2006; Hodyss et al. 2009; Martín-Doménech et al. 2016) are normally somewhat higher; around 1015 cm<sup>-1</sup> when CH<sub>3</sub>OH is mixed with H<sub>2</sub>O. Even though this band is rather weak, it offers a tool to unambiguously link to CH<sub>3</sub>OH as it has no overlap with bands of other species formed in these experiments and actually provides a nicely isolated feature.

A series of additional experiments are performed in order to further prove that the band at 1001 cm<sup>-1</sup> is indeed due to the C–O stretching frequency of CH<sub>3</sub>OH. This is realized in a set of isotope substitution experiments. Figure 4 shows spectra of three experiments that involve CH<sub>4</sub>, <sup>13</sup>CH<sub>4</sub>, O<sub>2</sub>, and <sup>18</sup>O<sub>2</sub> isotopes. By the simple harmonic oscillator approach, a heavier isotope should result in a redshift in the stretching frequency. As seen in Fig. 4, the 1001 cm<sup>-1</sup> feature redshifts to 984 cm<sup>-1</sup> and 973 cm<sup>-1</sup> for the reactions involving <sup>13</sup>CH<sub>4</sub> and <sup>18</sup>O<sub>2</sub>, respectively. This leads to differences of 17 cm<sup>-1</sup> and 28 cm<sup>-1</sup> from the 1001 cm<sup>-1</sup> feature, respectively, comparable to the differences found in CH<sub>3</sub>OH isotope experiments performed by Maity et al. (2014). These findings are consistent with isotopically enriched C–O bonds in newly formed CH<sub>3</sub>OH.

### 3.1.3. Temperature-dependent RAIR difference spectra

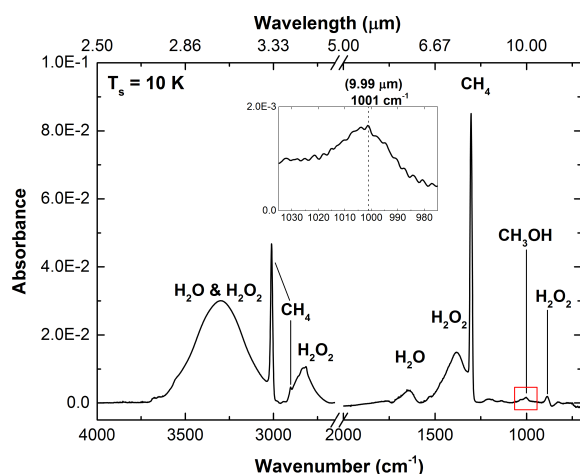
The TPD and RAIRS data can be correlated to one another to match IR features with desorption temperatures to further identify species that are initially made in the ice (i.e., before thermal processing via the TPD technique). Figure 5 shows the change in the IR features as a function of the substrate temperature (i.e., RAIR difference spectra) for the CH<sub>4</sub> + H +

**Table 2.** Infrared ice signatures from the  $\text{CH}_4 + \text{H} + \text{O}_2$  reaction and the corresponding vibrational modes.

Peak position ( $\text{cm}^{-1}$ )	Peak position ( $\mu\text{m}$ )	Molecule	Mode*	Reference
884	11.3	$\text{H}_2\text{O}_2$	$\nu_3$	1, 2, 3
1001	9.99	$\text{CH}_3\text{OH}$	$\nu_8$	this work
1302	7.68	$\text{CH}_4$	$\nu_4$	4, 5, 7, 8, 9
1381	7.24	$\text{H}_2\text{O}_2$	$\nu_2$	1, 2, 3
1637	6.11	$\text{H}_2\text{O}$	$\nu_2$	6
2815	3.55	$\text{CH}_4$	$\nu_2 + \nu_4$	4, 5, 7, 8, 9
2836	3.53	$\text{H}_2\text{O}_2$	$2\nu_2$	1, 2, 3
2902	3.45	$\text{CH}_4$	$\nu_1$	7, 8, 9
3010	3.32	$\text{CH}_4$	$\nu_3$	4, 5, 7, 8, 9
3295	3.0	$\text{H}_2\text{O}_2 + \text{H}_2\text{O}$	O–H stretch	10, 11
3675	2.7	$\text{H}_2\text{O}$	dangling bonds	7

**Notes.** (\*) The vibrational mode numbers are obtained from NIST.

**References** [1] Giguere & Harvey (1959), [2] Lannon et al. (1971), [3] Romanzin et al. (2011), [4] Hagen et al. (1981), [5] Gerakines et al. (2005), [6] Hodyss et al. (2009), [7] Gálvez et al. (2009), [8] Herrero et al. (2010), [9] Ennis et al. (2011), [10] Cuppen et al. (2010), [11] Ioppolo et al. (2010).

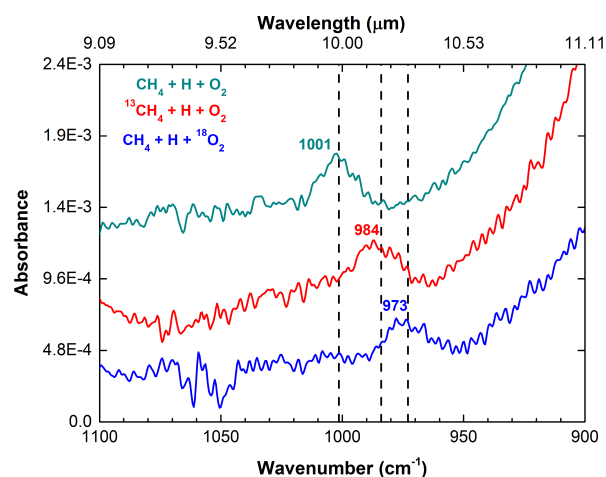


**Fig. 3.** RAIR spectrum of the reaction,  $\text{CH}_4 + \text{H} + \text{O}_2$ , on a 10 K substrate surface after 12 hours (exp. 1.0). The relative column density of  $\text{CH}_4:\text{CH}_3\text{OH}$  is 100:1, and the relative column density of  $\text{H}_2\text{O}:\text{CH}_3\text{OH}$  is 100:4. Spectra of the resulting species are shown (see Table 2), most noticeably a weaker band around  $1000\text{ cm}^{-1}$  (inset) highlighting the C–O stretch of the  $\text{CH}_3\text{OH}$  feature.

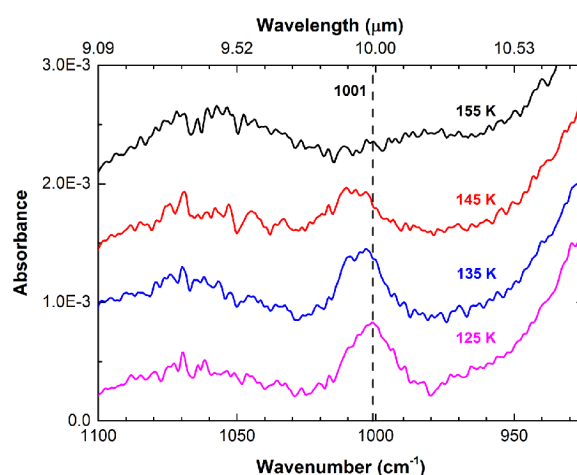
$\text{O}_2$  reaction. In Sect. 3.1.1 the desorption of  $\text{CH}_3\text{OH}$  is observed at 150 K, and in Sect. 3.1.2 the IR feature at  $1001\text{ cm}^{-1}$  is shown to be due to  $\text{CH}_3\text{OH}$  formation. The combination of these two pieces of data tells us that below 150 K,  $\text{CH}_3\text{OH}$  should be present in the ice, and above 150 K a majority of the  $\text{CH}_3\text{OH}$  should desorb. In Fig. 5, the C–O bond feature remains present at a substrate temperature of 145 K and is absent at a substrate temperature of 155 K. This correlates well with the desorption temperature of  $\text{CH}_3\text{OH}$  observed in Fig. 1 (bottom). Thus, the experiments discussed in Sects. 3.1.1–3.1.3 make a consistent case that  $\text{CH}_3\text{OH}$  is formed in the ice.

### 3.2. Spectral signature of $\text{CH}_3\text{OH}$ in a $\text{H}_2\text{O}$ -rich interstellar ice analogue

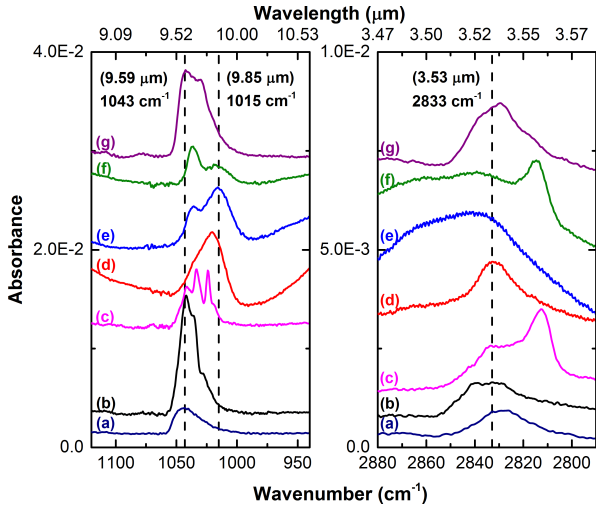
Because the C–O bond of  $\text{CH}_3\text{OH}$  is known to be sensitive to its surrounding environment, it holds a potential as a diagnostic tool to provide a deeper insight to the interactions of  $\text{CH}_3\text{OH}$  with



**Fig. 4.** RAIR spectra of the reactions,  $\text{CH}_4 + \text{H} + \text{O}_2$  (exp. 2.0),  $^{13}\text{CH}_4 + \text{H} + \text{O}_2$  (exp. 2.1), and  $\text{CH}_4 + \text{H} + ^{18}\text{O}_2$  (exp. 2.2) on a 10 K substrate surface after 6 hours each.



**Fig. 5.** RAIR difference spectra of the  $\text{CH}_4 + \text{H} + \text{O}_2$  reaction (exp. 2.0). The  $1001\text{ cm}^{-1}$  feature blueshifts between 125 K and 145 K, and could be due to  $\text{CH}_3\text{OH}$  segregating in a  $\text{H}_2\text{O}$ -ice environment (Martín-Doménech et al. 2014), and disappears upon further heating to 155 K.



**Fig. 6.** *Left:* RAIR spectra of the C–O stretch of CH<sub>3</sub>OH when CH<sub>3</sub>OH is mixed with different species. (a) Pure CH<sub>3</sub>OH (exp. 4.0), (b) CH<sub>3</sub>OH + O<sub>2</sub> (exp. 4.1), (c) CH<sub>3</sub>OH + CH<sub>4</sub> (exp. 4.2), (d) CH<sub>3</sub>OH + H<sub>2</sub>O (exp. 4.3), (e) CH<sub>3</sub>OH + H<sub>2</sub>O (exp. 4.4), (f) CH<sub>3</sub>OH + H<sub>2</sub>O + CH<sub>4</sub> (exp. 4.5), and (g) CH<sub>3</sub>OH + CO (exp. 4.6). *Right:* RAIR spectra of the symmetric C–H stretch of CH<sub>3</sub>OH.

other species in the H<sub>2</sub>O-rich interstellar ice analogue. Figure 6 (left) shows a RAIR spectrum of the C–O stretching frequency of pure CH<sub>3</sub>OH (a) in comparison to CH<sub>3</sub>OH embedded in an environment originating from the CH<sub>4</sub> + H + O<sub>2</sub> reaction (f), in addition to CH<sub>3</sub>OH mixed with other species (b–e, g). Figure 6 (right) shows a RAIR spectrum of the C–H stretch of pure CH<sub>3</sub>OH; a mode that is a useful tracer of CH<sub>3</sub>OH ice and can be probed by observational facilities (e.g., JWST NIRSpec). The experiments are performed using both the previously described HABS and a microwave atom source (MWAS), which dissociates a fraction of the incoming molecules into fragments. The fragments can recombine with each other and form products in the ice that would be found in, for example, the CH<sub>4</sub> + H + O<sub>2</sub> reaction (OH, H, O<sub>2</sub>, O, H<sub>2</sub>, H<sub>2</sub>O<sub>2</sub>, etc.). The C–O stretching frequency of pure CH<sub>3</sub>OH is found at 1043 cm<sup>-1</sup> and redshifts to 1015 cm<sup>-1</sup> in the ice where CH<sub>4</sub> and fragments from dissociated H<sub>2</sub>O are co-deposited. We note that the peak does not redshift to 1001 cm<sup>-1</sup>, and there is a logical explanation for this. As illustrated in Fig. 7 in the study by Dawes et al. (2016), an increase in the H<sub>2</sub>O concentration in a CH<sub>3</sub>OH + H<sub>2</sub>O mixture does not necessarily correlate to a more redshifted C–O stretching frequency. Thus, it is expected that a specific ratio of the ice products in the CH<sub>4</sub> + H + O<sub>2</sub> reaction is needed in order to recreate the 1001 cm<sup>-1</sup> feature. From the spectra presented in Fig. 6, it is clear that the largest redshift for the C–O stretch takes place in a H<sub>2</sub>O-rich ice matrix that contains CH<sub>4</sub>, H<sub>2</sub>O, and H<sub>2</sub>O dissociation products (i.e., the mixture in the actual CH<sub>4</sub> + H + O<sub>2</sub> experiment; exp. 4.5), as expected. When CH<sub>3</sub>OH is mixed with a single species, such as H<sub>2</sub>O or CH<sub>4</sub>, the C–O stretching frequency does not redshift as much. In conclusion, it is proposed that many, rather than one or two molecular species, interact simultaneously with CH<sub>3</sub>OH in the CH<sub>4</sub> + H + O<sub>2</sub> reaction.

### 3.3. CH<sub>3</sub>OH formation at 10 and 20 K

As mentioned in Sect. 2, an absolute value for the CH<sub>3</sub>OH column density is not obtained in this study. Thus, in order to provide information on the formation yield of CH<sub>3</sub>OH formed at 10–20 K, a relative yield is obtained by comparing the amount

of CH<sub>3</sub>OH formed at 10 K (exp. 2.0) to the amount of CH<sub>3</sub>OH formed at 20 K (exp. 5.0). Since it is expected in interstellar space that CH<sub>4</sub> and H<sub>2</sub>O ice are initially formed by hydrogenation reactions (Tielens & Hagen 1982; Öberg et al. 2008; Ioppolo et al. 2008; Cuppen et al. 2010; van Dishoeck et al. 2013), the CH<sub>4</sub>:H<sub>2</sub>O ice ratio should not be drastically different between surface temperatures of 10 and 20 K. Therefore, the CH<sub>4</sub>:H<sub>2</sub>O ice ratio is kept the same in this temperature range. At 20 K, a 65% decrease in the CH<sub>3</sub>OH abundance is observed compared to the amount of CH<sub>3</sub>OH formed at 10 K. This decrease in abundance is expected to be due to the decrease in the H-atom lifetime on the surface (Fuchs et al. 2009; Cuppen et al. 2010), as expected if CH<sub>3</sub>OH is being formed by the OH radical. Although, it should be noted that the reaction steps for CH<sub>3</sub>OH formation may also be affected by the temperature change, which can influence the change in the CH<sub>3</sub>OH abundance.

### 3.4. Constraining the formation of CH<sub>3</sub>OH in the overall reaction network

The experiments discussed here show that it is possible to form CH<sub>3</sub>OH in a H<sub>2</sub>O-rich environment along the low temperature pathway:



To observe how these channels fit into the overall reaction network ongoing in H<sub>2</sub>O-rich ice, the change in abundances of CH<sub>3</sub>OH and other ice species needs to be tracked (Öberg et al. 2010). Figure 7 shows two RAIR spectra of the reaction, CH<sub>4</sub> + H + O<sub>2</sub>, and the subsequent infrared features pertaining to newly formed CH<sub>3</sub>OH, H<sub>2</sub>O, H<sub>2</sub>O<sub>2</sub>, and O<sub>3</sub>. With an enhancement of the O<sub>2</sub> flux (dashed line in Fig. 7), the CH<sub>3</sub>OH abundance decreases and the H<sub>2</sub>O<sub>2</sub> and O<sub>3</sub> abundances clearly increase. The decrease in the CH<sub>3</sub>OH abundance at the expense of the H<sub>2</sub>O<sub>2</sub> and O<sub>3</sub> abundances gives an indication that the formation channel of CH<sub>3</sub>OH at least crosses with the formation channels of H<sub>2</sub>O<sub>2</sub> and O<sub>3</sub>.

The link between the formation of CH<sub>3</sub>OH, H<sub>2</sub>O<sub>2</sub>, and O<sub>3</sub> in the H<sub>2</sub>O-rich ice can be better understood when looking at the reaction network presented by Cuppen et al. (2010). In this study, the H + O<sub>2</sub> reaction is extensively discussed, and it is found that the O atom used to react with O<sub>2</sub> to form O<sub>3</sub> is most likely created by two reactions,



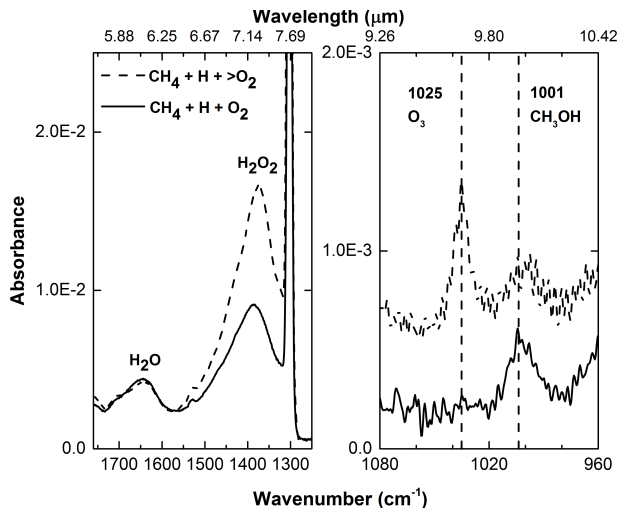
and



where OH is mainly formed by the following reaction:



We note that only the relevant steps are discussed here, and the full reaction scheme can be found in Cuppen et al. (2010). The RAIR spectra shown in Fig. 7 suggest that the latter reactions are dominant in the CH<sub>4</sub> + H + O<sub>2</sub> reaction. As observed in Fig. 7, a threefold increase in the O<sub>2</sub> flux (exp. 2.0 has a threefold increase in the O<sub>2</sub> flux compared to exp. 6.0) results in a substantial increase in the H<sub>2</sub>O<sub>2</sub> abundance compared to the change in the H<sub>2</sub>O abundance. This is consistent with reactions (3a) and (3b) as HO<sub>2</sub> is an intermediate of the

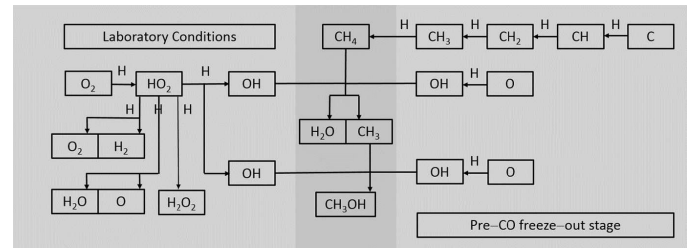


**Fig. 7.** RAIR spectra of the reactions,  $\text{CH}_4 + \text{H} + \text{O}_2$ , where  $>\text{O}_2$  indicates a threefold increase in the  $\text{O}_2$  flux. The dashed-line spectrum corresponds to exp. 2.0, and the solid-line spectrum corresponds to exp. 6.0.

$\text{H}_2\text{O}_2$  product;  $\text{HO}_2$  is also an intermediate of  $\text{H}_2\text{O}$  formation, as shown in reaction (2). However, comparison of the two spectra in Fig. 7 shows that an increase in the  $\text{O}_3$  abundance is followed by a significant difference in the  $\text{H}_2\text{O}_2$  abundance and a relatively minor change in the  $\text{H}_2\text{O}$  abundance, suggesting that reactions (3a) and (3b) are carried out more efficiently than reaction (2) in the  $\text{CH}_4 + \text{H} + \text{O}_2$  experiment. Decreasing the  $\text{CH}_4$  abundance by 20% shows a 10% decrease in the  $\text{CH}_3\text{OH}$  abundance and no change in the  $\text{H}_2\text{O}$  abundance, which further supports that reaction (2) is not the dominating channel. In addition, reactions (3a) and (3b) show that  $\text{O}_3$  is formed by an OH-induced process, and this is expected since  $\text{CH}_3\text{OH}$  is formed by an OH-mediated reaction. These insights piece together to demonstrate that reactions (3a) and (3b) are the likely pathways that connect the formation channel of  $\text{CH}_3\text{OH}$  to other formation channels in a  $\text{H}_2\text{O}$ -rich ice reaction network. In Fig. 8, the laboratory experiments and the astrochemical proposed reactions are summarized in one figure. The main difference between the two is how OH radicals are inserted into the reaction scheme, but this does not change the conclusions derived here for the presented pre-CO freeze-out formation scheme.

#### 4. Astrophysical implications

The experimental work presented here calls into question the  $A_V$  at which  $\text{CH}_3\text{OH}$  ice starts to grow (i.e., the  $\text{CH}_3\text{OH}$  ice formation threshold). If an observational search for the  $\text{CH}_3\text{OH}$  ice formation threshold in low extinction (i.e.,  $\text{H}_2\text{O}$ -rich) environments is to be performed, it would be useful to know beforehand an approximation of the  $\text{CH}_3\text{OH}$  abundance in these  $\text{H}_2\text{O}$ -rich ices. That way, the amount of  $\text{CH}_3\text{OH}$  formed before and after the CO freeze-out stage can be distinguished. A laboratory experiment of  $\text{CO} + \text{H}$  performed under similar conditions to exp. 2.0 shows that the  $\text{CO}:\text{CH}_3\text{OH}$  ratio is 100:20. In the parallel  $\text{CH}_4 + \text{OH}$  experiment, the  $\text{CH}_4:\text{CH}_3\text{OH}$  ratio is 100:1. A comparison of these two experiments leads to the conclusion that the  $\text{CH}_4 + \text{OH}$  reaction is 20 times less efficient at producing  $\text{CH}_3\text{OH}$  than the  $\text{CO} + \text{H}$  reaction in the laboratory setting. Making the simple assumption that  $\text{CH}_3\text{OH}$  is primarily formed by  $\text{CH}_4 + \text{OH}$  and  $\text{CO} + \text{H}$  in interstellar ices around the CO freeze-out stage, laboratory data can be combined with

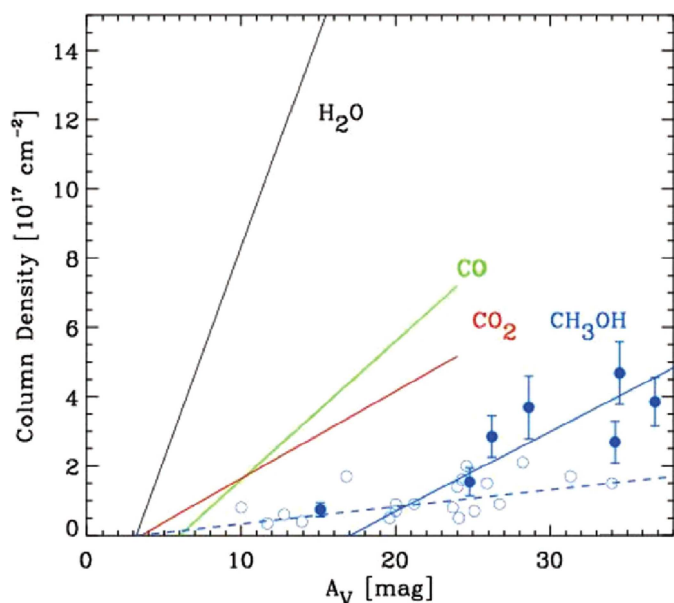


**Fig. 8.** Comparison of the laboratory work discussed here and the astrochemical network, resulting in the formation of  $\text{CH}_3\text{OH}$  in the pre-CO freeze-out stage. Only the relevant steps are shown.

observational data to determine the amount of  $\text{CH}_3\text{OH}$  in the two different ice phases. According to Boogert et al. (2015), the  $\text{H}_2\text{O}:\text{CH}_4$  and  $\text{H}_2\text{O}:\text{CO}$  median ice ratios around low mass young stellar objects (LYSOs) are 100:4.5 and 100:21, respectively. Incorporation of the laboratory results shows that 0.045% of the interstellar ice should contain  $\text{CH}_3\text{OH}$  that derives from  $\text{CH}_4$ , and 4.2% of the ice should contain  $\text{CH}_3\text{OH}$  that derives from CO. Thus, of the total observed abundance of  $\text{CH}_3\text{OH}$ , about 1% of the total amount is the abundance of  $\text{CH}_3\text{OH}$  from the  $\text{CH}_4 + \text{OH}$  reaction that is expected to be ongoing in the  $\text{H}_2\text{O}$ -rich ice phase. Although a  $\text{CH}_3\text{OH}:\text{H}_2\text{O}$  percentage of 0.045% may appear insignificant, multiplying this value by the typical  $\text{H}_2\text{O}$  ice to H ratio of  $2\text{E}-5$  yields a  $\text{CH}_3\text{OH}$  ice to H ratio of  $9\text{E}-9$ , which is greater than the abundance of gas-phase  $\text{CH}_3\text{OH}$  and other gas-phase species detected in cold cores relative to  $\text{H}_2$  (Herbst & Van Dishoeck 2009). Using values from the 20 K experiment (exp. 5.0), the  $\text{CH}_3\text{OH}$  ice to H ratio results in a value of  $3\text{E}-9$ , which is similar to the abundance of gas-phase  $\text{CH}_3\text{OH}$  in these particular astrophysical environments. We note that the median ice ratios contain sources that also have a relatively high amount of CO ice, so the calculated values are not entirely representative of low extinction (i.e.,  $\text{H}_2\text{O}$ -rich ice) scenarios. See Fig. 9 for a better approximation of the possible  $\text{CH}_3\text{OH}$  abundance in the  $\text{H}_2\text{O}$ -rich ice layer at various extinctions.

Figure 9, which is adapted from Boogert et al. (2015), shows that  $\text{CH}_3\text{OH}$  detections have an ice formation threshold at an  $A_V$  value higher than  $\text{H}_2\text{O}$  and CO (the ice formation threshold in this figure is the X-axis cut-off divided by 2 because the background stars trace both the front and back sides of the clouds). This is in line with the currently accepted model that  $\text{CH}_3\text{OH}$  ice is mainly formed in prestellar cores at  $A_V > 9$  via hydrogenation of accreted CO molecules. However, what is puzzling is that the upper limit sight lines are all below the detections and at similar  $A_V$  values to that of the detections. The  $\text{CH}_3\text{OH}:\text{H}_2\text{O}$  ratio of 4:100 derived from the laboratory work presented here (from exp. 1.0) is represented by the dashed blue line and is at the level of  $3\text{-}\sigma$  upper limits. Thus, the current observations are consistent with a  $\text{CH}_3\text{OH}:\text{H}_2\text{O}$  ratio of at most a few percent, though likely not as high as 4%, the value seen in our optimized experiments, since  $\text{H}_2\text{O}$  is a by-product in the  $\text{CH}_4 + \text{H} + \text{O}_2$  reaction and is not the main source of  $\text{H}_2\text{O}$  in interstellar ices. At this level, a formation threshold as low as that of  $\text{H}_2\text{O}$  (1.6 mag) as a result of chemical reactions in  $\text{H}_2\text{O}$ -rich ices cannot be excluded. More sensitive observations are needed to further constrain this scenario and with the upcoming launch of the JWST, this soon will be within range.

The insights gained here also link to the formation of larger COMs. Reaction (1) gives rise to the possibility that such a mechanism could lead to the formation of more complex species in  $\text{H}_2\text{O}$ -rich ices. It has been shown that COMs such as



**Fig. 9.** Relationships between ice column density and visual extinction ( $A_V$ ) as observed toward stars behind dense clouds and cores for  $\text{H}_2\text{O}$  (black),  $\text{CO}$  (green),  $\text{CO}_2$  (red), and  $\text{CH}_3\text{OH}$  (blue). For clarity, for  $\text{H}_2\text{O}$ ,  $\text{CO}$ , and  $\text{CO}_2$ , the least-squares linear fits are shown and not the actual data points. For  $\text{CH}_3\text{OH}$ , blue bullets indicate detections with  $1-\sigma$  error bars and open circles indicate  $3-\sigma$  upper limits. The solid blue line is a fit to the detections, and the dashed blue line (from exp. 1.0) indicates the scenario in which  $\text{CH}_3\text{OH}$  has the same ice formation threshold as  $\text{H}_2\text{O}$  at a maximum abundance of  $\text{CH}_3\text{OH}:\text{H}_2\text{O} = 4\%$ , as found in the listed experiments. The  $3\sigma$  upper limits prove in principle that  $\text{CH}_3\text{OH}:\text{H}_2\text{O}$  is less than 4% in the observed sight lines. This figure was adapted from Fig. 7 in Boogert et al. (2015).

glycolaldehyde, ethylene glycol, and glycerol can be successfully formed by hydrogenation of  $\text{CO}$  (Fedoseev et al. 2015, 2017), thus providing a promising reaction scheme for COM formation in  $\text{CO}$ -rich interstellar ices. Formation of COMs in  $\text{H}_2\text{O}$ -rich ices has been less studied, and the mechanism described in this paper could be a potential source of molecular complexity in cold  $\text{H}_2\text{O}$ -dominated ice matrices, although formation efficiencies will be lower.

## 5. Conclusions

The formation of  $\text{CH}_3\text{OH}$  by  $\text{CH}_4 + \text{OH}$  shown in this laboratory study suggests that  $\text{CH}_3\text{OH}$  ice can also be formed before the heavy  $\text{CO}$  freeze-out stage in prestellar cores, i.e., at  $A_V < 3$  and grain temperatures of 20 K. The main findings from this work are summarized below:

- The formation of  $\text{CH}_3\text{OH}$  in  $\text{H}_2\text{O}$ -rich ices occurs by the sequential reactions,  $\text{CH}_4 + \text{OH} \rightarrow \text{CH}_3 + \text{H}_2\text{O}$  and  $\text{CH}_3 + \text{OH} \rightarrow \text{CH}_3\text{OH}$ .
- Since  $\text{CH}_4$  and  $\text{OH}$  radicals are expected to be found in the  $\text{H}_2\text{O}$ -rich ice phase of cold dense cores,  $\text{CH}_3\text{OH}$  is also expected to be found in those ices, suggesting that the  $\text{CH}_3\text{OH}$  formation threshold is below  $A_V = 9$ .
- Raising the sample temperature to 20 K results in a 65% decrease in the  $\text{CH}_3\text{OH}$  abundance compared to  $\text{CH}_3\text{OH}$  formed on a 10 K surface, showing that the formation pathway of  $\text{CH}_3\text{OH}$  in this study is relevant to the period of the cold dense core stage at which ice species are formed primarily by atom-induced reactions (i.e., before gas phase molecules, like  $\text{CO}$ , accrete onto the grain surface).

- Under similar experimental parameters, the  $\text{CO} + \text{H}$  channel is about 20 times more efficient to forming  $\text{CH}_3\text{OH}$  than the  $\text{CH}_4 + \text{OH}$  channel at temperatures around 10 K.
- More sensitive astronomical observations are warranted to determine the  $\text{CH}_3\text{OH}:\text{H}_2\text{O}$  ratio in  $\text{H}_2\text{O}$ -rich interstellar ices, which is likely to be a few percent at most, as derived from Fig. 9.
- The formation of COMs under cold dark cloud conditions is mainly linked to the  $\text{CO}$  freeze-out stage. The reactions studied here carry the potential to lead to the formation of COMs at an earlier astrochemical evolution stage.

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